

# PREDICTIONS IN THE LANDSCAPE

RAPHAEL BOUSSO, UC BERKELEY

IN MEMORY OF ANDREW CHAMBLIN

OCTOBER 14, 2006

CAMBRIDGE

# COLLABORATORS

hep-th /

A. Chamblin

9805167, 0004134

J. Polchinski

B. Freivogel

0603105, 0606114

I. Yang

0605263, 0610132

R. Harnik

G. Kribs

G. Perez

Holographic



Entropic

LANDSCAPE

I. Survey long-lived vacua

II. Probabilities in Eternal Inflation

III. Anthropic selection



PREDICTIONS

**Rees**

Weinberg

Sakharov

Linde

GARRIGA

Page

Lin

VILENKO

Dyson

Kleban

Hertog

Susskind

Banks

Hawking

Winitski

Vancharin

Mehlumian

## Easter

Martin

Schwartz-Perlov

Tegmark

Hall

Aguirre

## Feldstein

Watari

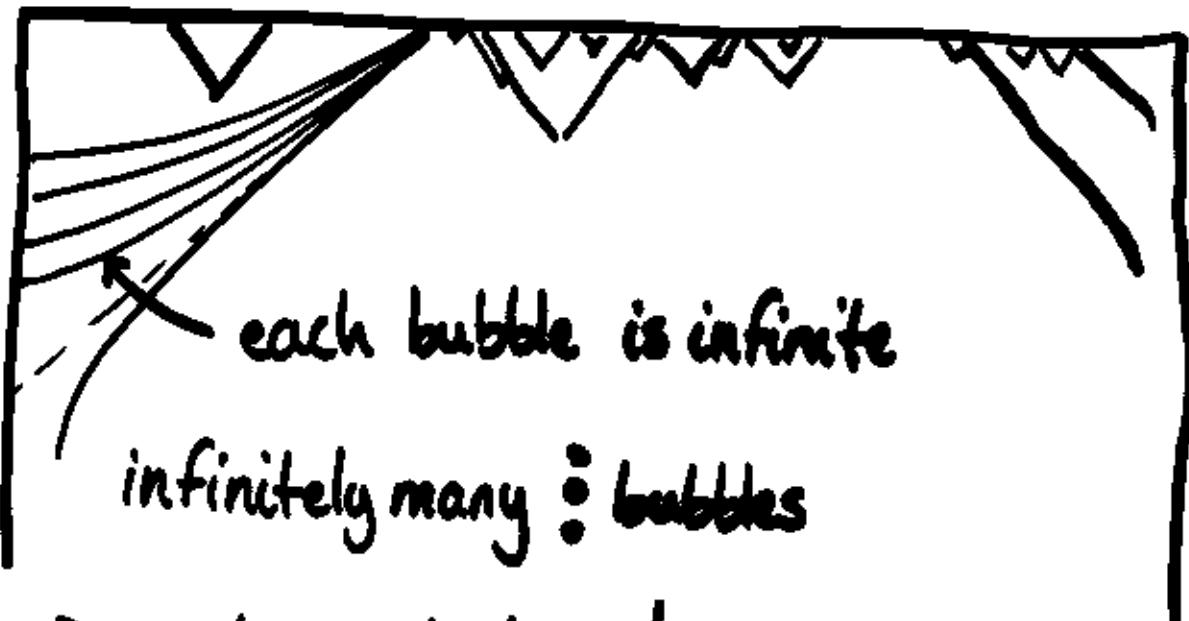
Gibbons

Turck

Johnson

## II.

# Eternal Inflation: Global Structure



---

At finite time, compare...

Garriga & Vilenkin  
gr-qc/0102090

Garriga et al.  
hep-th/0509184

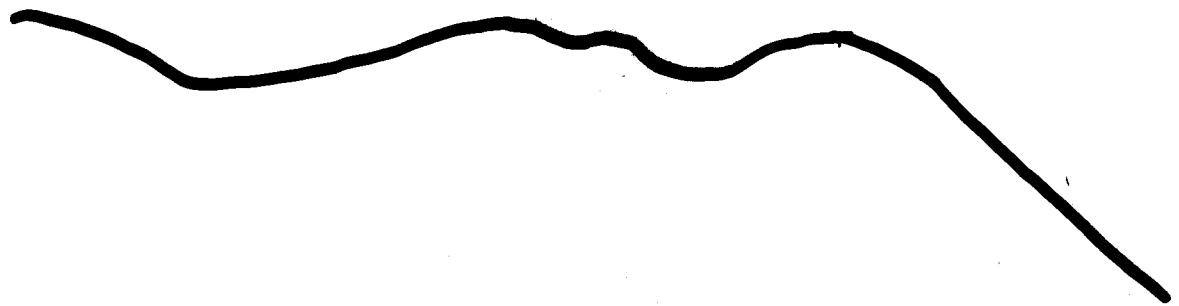
Easther et al.  
astro-ph/0511233

total volume of  $\ddot{\circ}$  bubbles

number of  $\ddot{\circ}$  bubbles

...

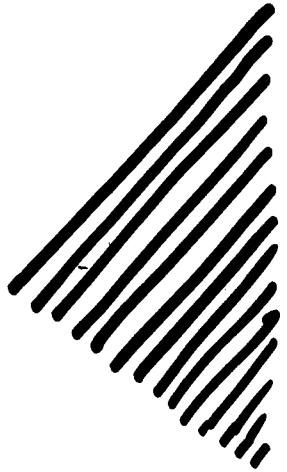
?



No preferred global time

→ get any answer you want !

[Linde et al., gr-qc/9601005]



Global ("bird's eye") view leads to  
ambiguities and pathologies.

Only one causally connected region is accessible  
to an observer.

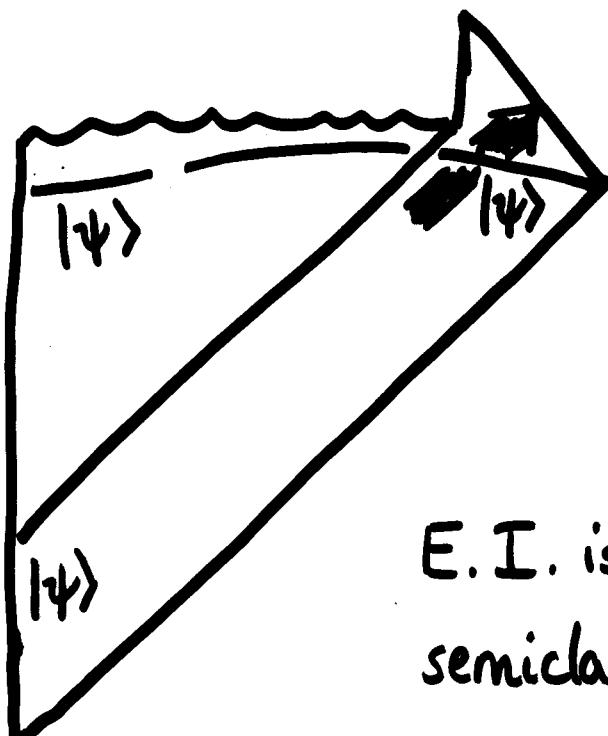
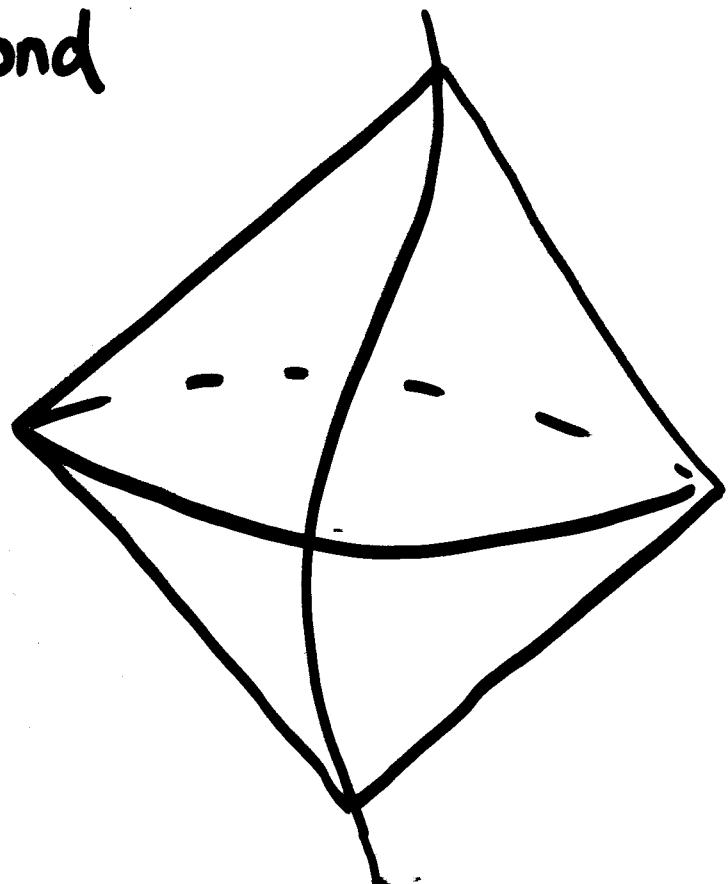
Along any generic worldline, inflation eventually ends.\*)

→ use Causal Diamond

as a regulator to  
define probabilities

Also motivated by

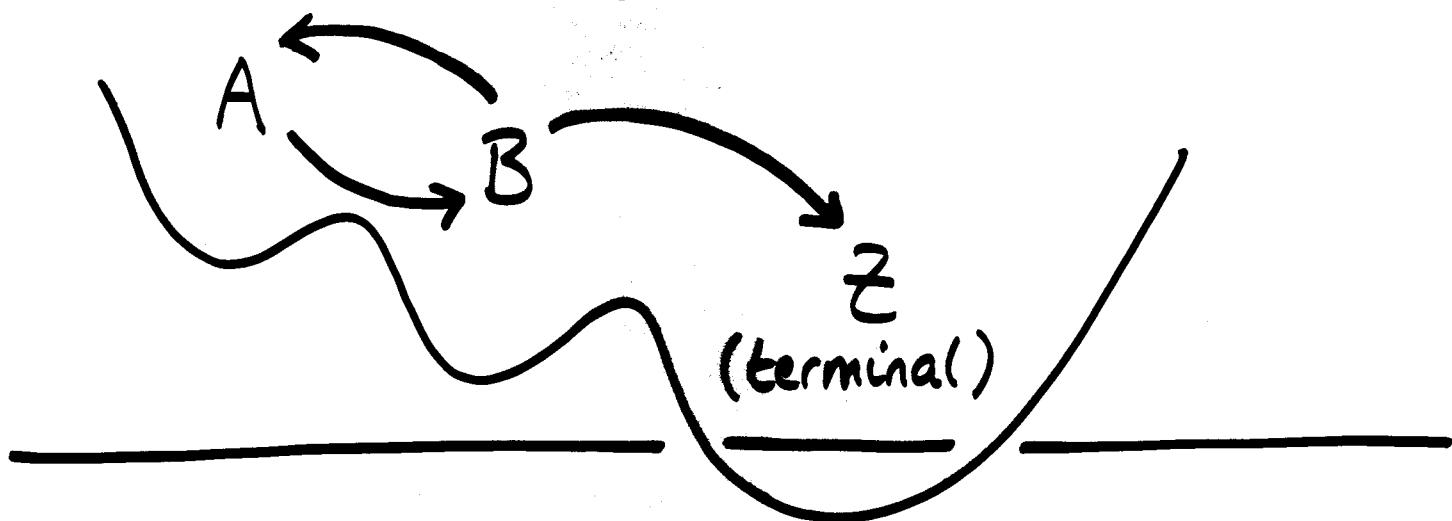
- Occam
- Unitarity in black hole evaporation



E.I. is worse : no unique  
semiclassical geometry outside  
causal diamond

Use a single worldline to compute probabilities

Consider a landscape,



start in A.

$K_{ij}$  = probability per unit time for  
worldline in vacuum  $j$  to  
enter vacuum  $i$

$$= \begin{pmatrix} 0 & * & 0 \\ * & 0 & 0 \\ 0 & * & 0 \end{pmatrix}$$

---

At finite time, compare...

Garriga & Vilenkin  
gr-qc/0102090

Garriga et al.  
hep-th/0509184

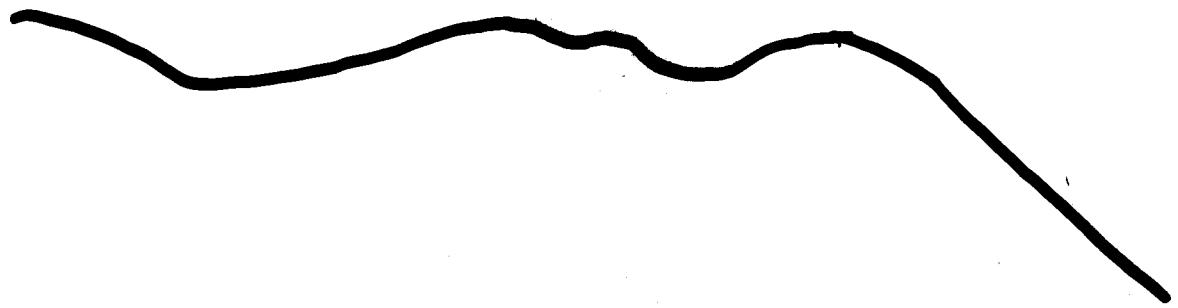
Easther et al.  
astro-ph/0511233

total volume of  $\ddot{\circ}$  bubbles

number of  $\ddot{\circ}$  bubbles

...

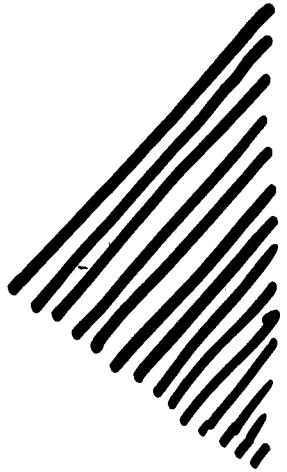
?



No preferred global time

→ get any answer you want !

[Linde et al., gr-qc/9601005]



Global ("bird's eye") view leads to  
ambiguities and pathologies.

Only one causally connected region is accessible  
to an observer.

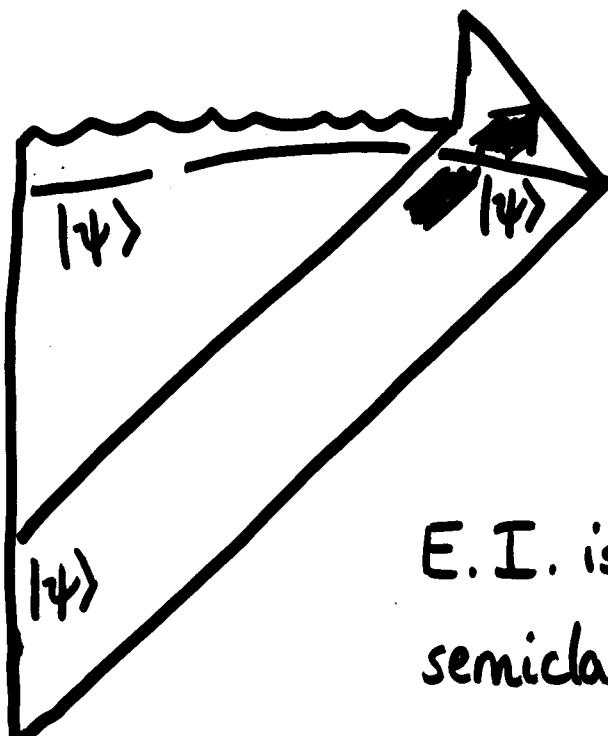
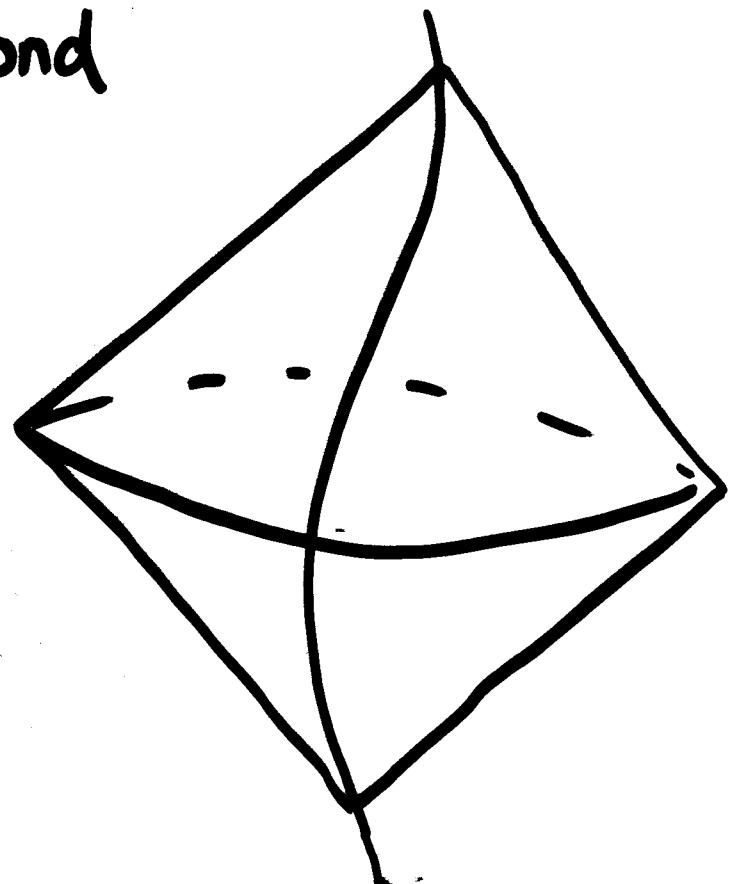
Along any generic worldline, inflation eventually ends.\*)

→ use Causal Diamond

as a regulator to  
define probabilities

Also motivated by

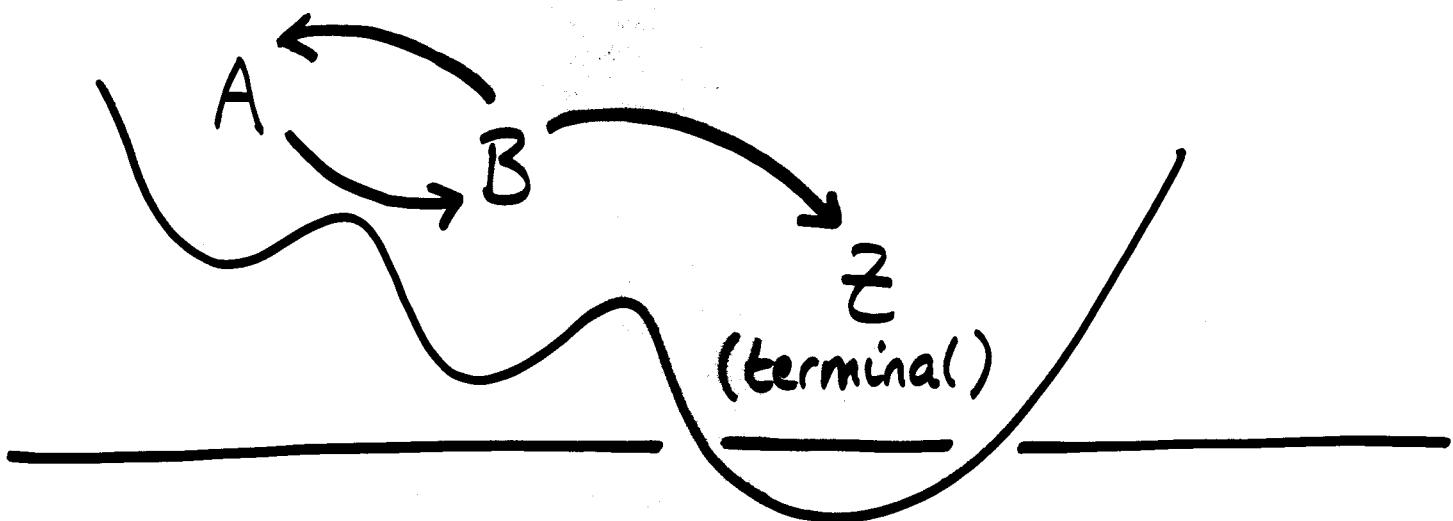
- Occam
- Unitarity in black hole evaporation



E.I. is worse : no unique  
semiclassical geometry outside  
causal diamond

Use a single worldline to compute probabilities

Consider a landscape,



start in A.

$K_{ij}$  = probability per unit time for  
worldline in vacuum j to  
enter vacuum i

$$= \begin{pmatrix} 0 & * & 0 \\ * & 0 & 0 \\ 0 & * & 0 \end{pmatrix}$$

### III.

Current approaches tend to break up the probability for vacuum  $i$  to be observed:

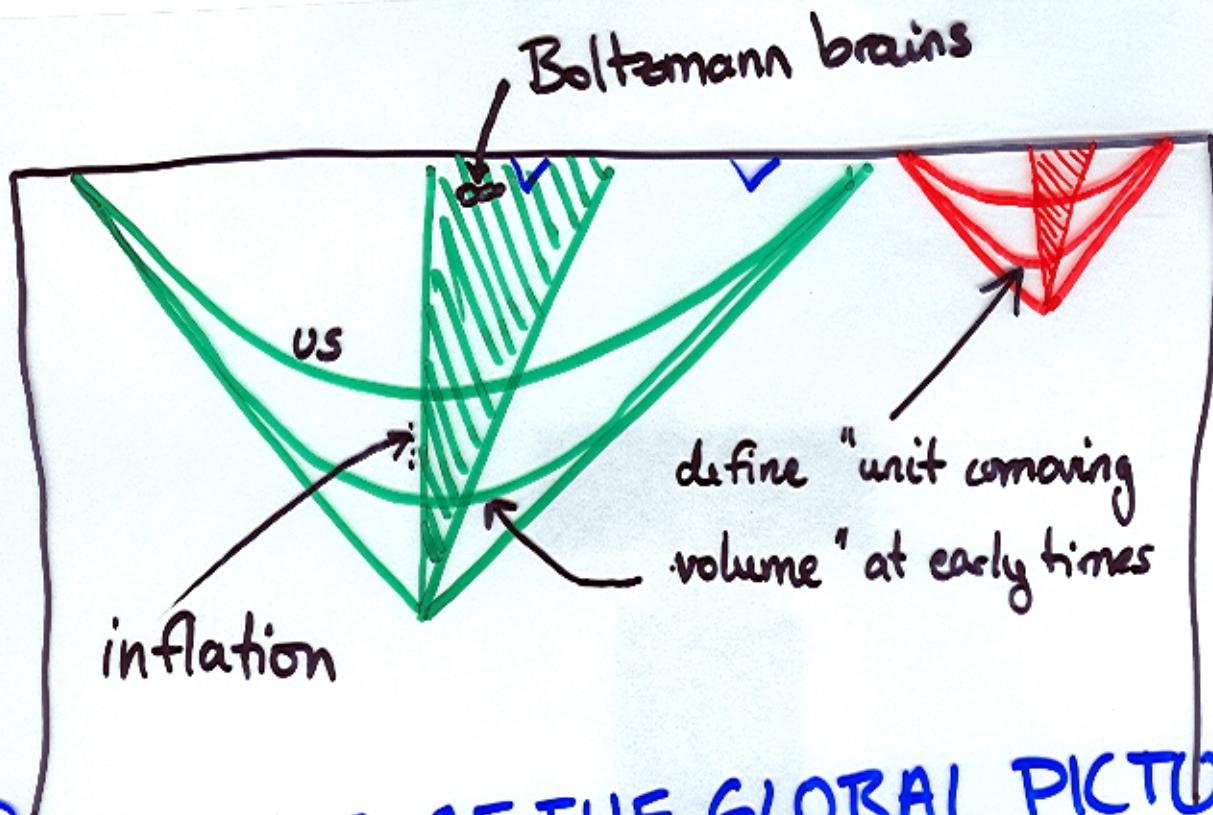
$$\Pi_i = p_i \omega_i$$

$p_i$ : probability for  $i$ -bubble to be produced

$\omega_i$ : usually taken to be proportional to the number of observers expected in the  $i$ -bubble. This is infinite in the global picture (or zero).  $\rightarrow$

Define  $\omega_i = \frac{N_{\text{obs}}}{\text{unit comoving volume}}$

(Vilenkin)



## PROBLEMS OF THE GLOBAL PICTURE:

- 1) How to define "observer" and estimate their formation rate in unfamiliar vacua?
- 2)  $\omega_i$  is exponentially sensitive to the duration of slow-roll inflation in the  $i$ -bubble  
 $\rightarrow \text{expect } \frac{\delta g}{g} \rightarrow \begin{cases} 0 & [\text{Feldstein, Watari} \\ & \& \text{Hall}], [\text{Vilenkin}] \\ 1 & \text{or} \end{cases}$
- 3) The overwhelming majority of observers are not like us. [Page] [RB, Freivogel]

The holographic cutoff makes vol(i)

finite



and admits

a weighting which mitigates ~~problems~~

problem #1. Solves #2 (see later), #3\*.

Observers require free energy

Must be able to increase entropy

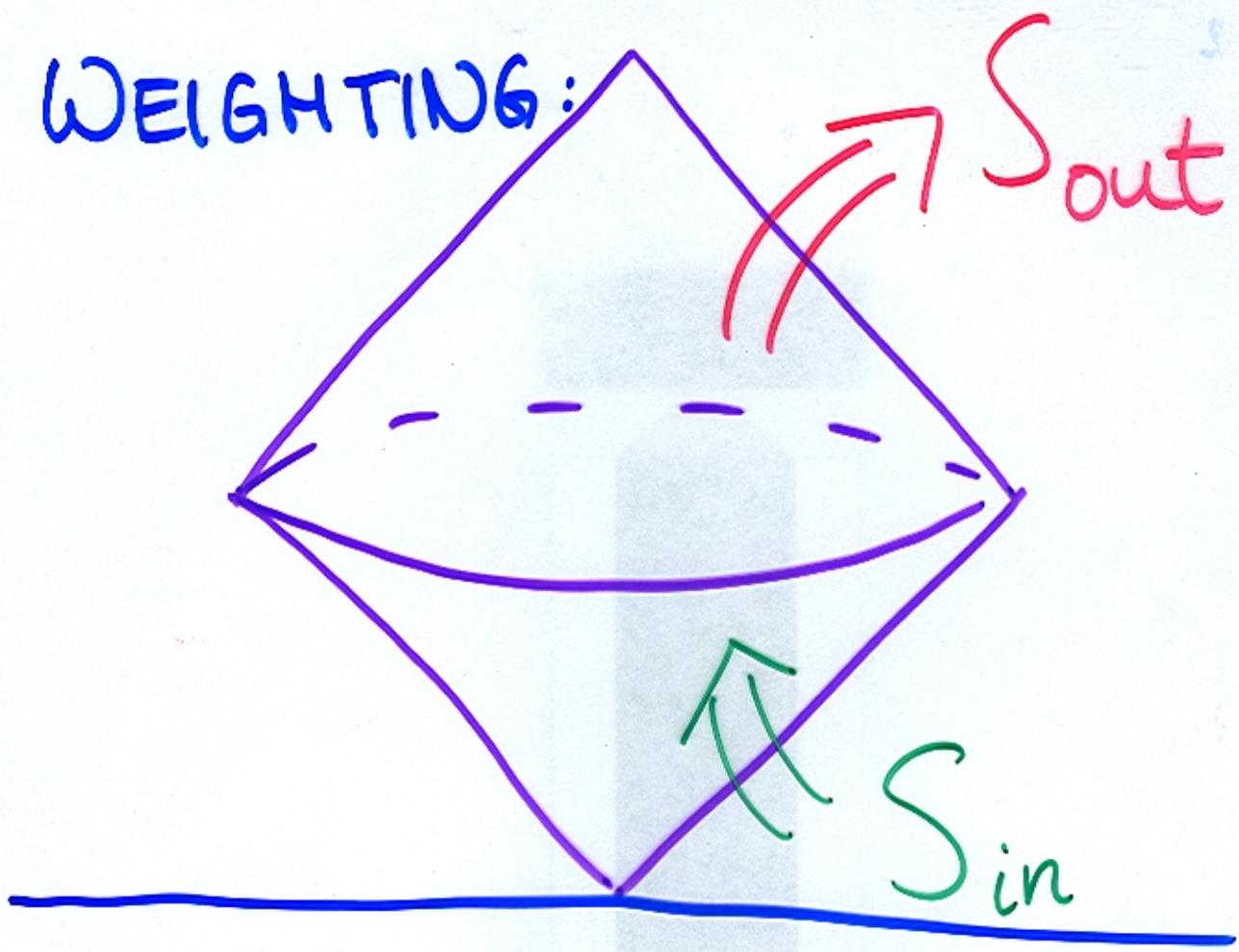
Estimate potential complexity of a vacuum by how much it allows

the entropy to increase within one

causal diamond :  $\Delta S \sim \frac{F}{T}$

Expect this to capture, e.g., structure formation.

# THE NOT-SO-ANTHROPIC WEIGHTING:



Define

$$\omega_i = \Delta S \equiv S_{\text{out}} - S_{\text{in}}$$

What  $\Delta S$  will not depend on :

- lifetime of vacuum ( $beyond \Lambda^{-\frac{1}{2}}$ )
- inflationary volume expansion ( $beyond$  suppressing curvature domination until  $\Lambda$  dominates).

Thus the local viewpoint resolves a paradox identified by Feldstein, Hall and Watari : If # of e-folds did enter the weight (as in most proposals) then  $\frac{\delta g}{g}$  would be driven towards an extreme value (0 or 1).

So much for problems #3 and #2.

Before turning to #1 there is a worry.

I did not use explicit anthropic weighting requiring  $\omega_i \rightarrow 0$  if no observers, only used  $\omega_i = \Delta S$ . Won't we lose the successful pre-/post-dictions of the anthropic principle?

E.g.  $\Lambda$ ,  $\frac{\delta g}{g}$ , ...

Test this idea on our data point.

(Ignore horizon entropy.)

In our universe, the main contribution

to  $\Delta S$  since reheating comes from  
stellar burning !  $(10^5 n_b)$

This requires not only structure formation

but galaxy formation (cooling  $\rightarrow \Delta S$ )

and long-lived stars.

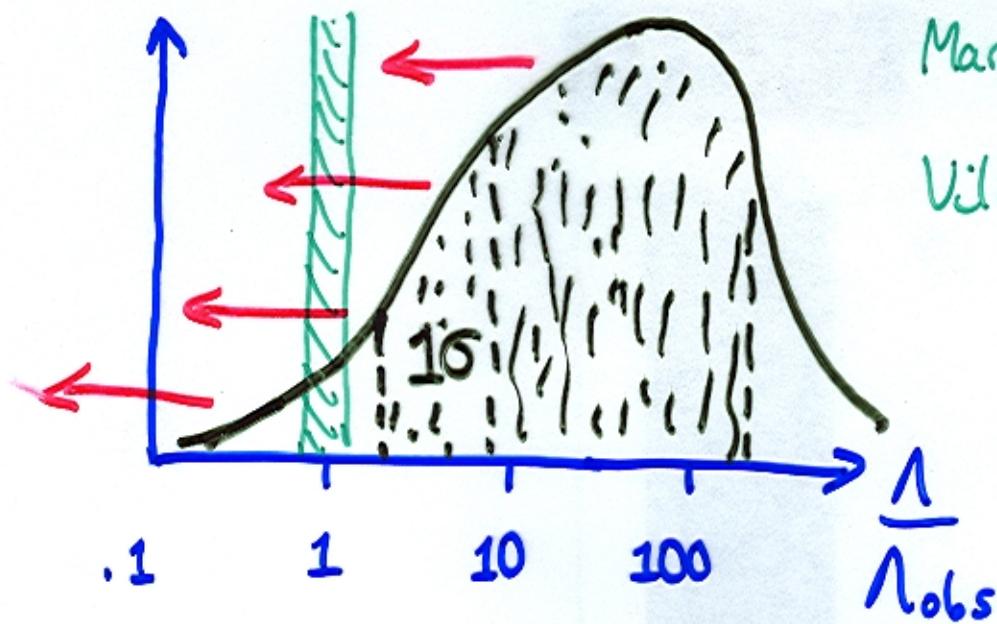
This result is a major success.

① This means that the entropic principle will reproduce the successful predictions of the anthropic principle, which were based on prior measurements:

- $t_\Lambda \gtrsim t_{\text{Galaxies}}$  or else  $\Delta S \rightarrow 0$
- $\frac{\delta S}{S} \gtrsim 10^{-6}$  or else  $\Delta S \rightarrow 0$
- coincidence problem: we should live at the time when most of  $\Delta S$  is produced:  
 $\sim 5 \text{ Gyr after } t_{\text{Galaxies}}$  ✓

Same results from much simpler assumption!

② In some cases, the entropic principle gives better agreement with observation



Martel Shapiro Weinberg  
Vilenkin et al.

Extra weight by total mass inside causal diamond,  $\Lambda^{-1/2}$ , favors smaller  $\Lambda$  and will shift curve to the left !

③ In our vacuum,  $\Delta S$  captures anthropic requirements usually put in by hand.

→ entropic weighting may ~~not~~ estimate (at least crudely) the observer content of very different vacua .

"Entropic Principle"  
(progress on problem #1.)

Can we predict  $\Lambda$  without priors?

(aka, What happens to the Weinberg bound if everything scans?)

$$p(\Lambda_1 < \Lambda < \Lambda_2) \propto \sum_i w_i ; w_i = \Delta S_i$$

need statistical averaging; define

$$W(\Lambda) \Delta \Lambda = \sum_i w_i ; N \ll \Delta \Lambda \ll 1 .$$

vacua in  $\Delta \Lambda$  interval

so that

$$p(\Lambda_1 < \Lambda < \Lambda_2) \propto \int_{\Lambda_1}^{\Lambda_2} d\Lambda W(\Lambda) .$$

# How to estimate $\omega(\lambda)$

$\omega(\lambda)$  is the average  $\Delta S$ .

Assumption: Suppose the average  $\Delta S$  scales with  $\lambda$  like the maximum  $\Delta S$ :

$$\omega(\lambda) = \alpha \Delta S_{\max}(\lambda)$$

We know that  $\Delta S_{\max} = 3\pi/\lambda$  (2<sup>nd</sup> law)

(This can actually be attained:

$$\Delta S \approx \frac{\text{energy burned}}{\text{temperature}} \lesssim \frac{\lambda^{-1/2}}{\lambda^{1/2}} = \frac{1}{\lambda} .$$

So  $\omega(\lambda)$  scales like  $\frac{1}{\lambda}$ .

$$\Rightarrow P(\Lambda_1 < \Lambda < \Lambda_2) = \int_{\Lambda_1}^{\Lambda_2} d\Lambda \omega(\Lambda)$$

$$\sim \boxed{\log \Lambda_2 - \log \Lambda_1}$$

Use structure of theory, e.g.  $N \approx 10^{500}$

→ prior-free prediction

$$\boxed{-\log \Lambda \sim 0(100)}$$

