

Structure formation in non-minimally coupled dark energy models

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Outline

- 1 **Motivations of the study**
- 2 **Models**
- 3 **Spherical Collapse Model**
- 4 **Observables**
- 5 **Conclusions**

Motivations

- Check the validity of N-body simulations
- Study the effect of a varying G on structure formation
- Compare models with the same $w(a)$ but with different physical background
- Compare models with the same $H(a)$ but with constant G
- Can we constrain them?

Models

- Λ CDM model (fiducial model)
- Minimally coupled models (to test the effects of $G(t)$)

$$S = \int d^4x \sqrt{-g} \left(\frac{R}{16\pi G} + \mathcal{L}_\phi + \mathcal{L}_\text{fl} \right)$$

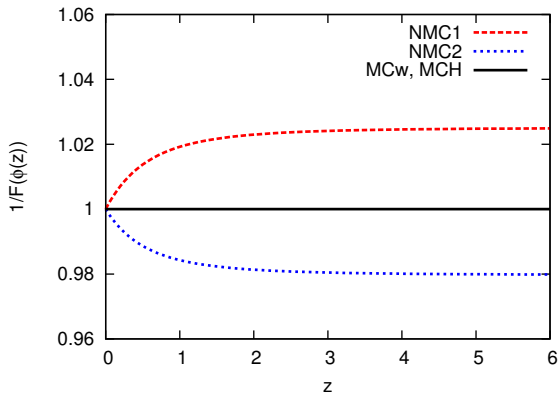
$$G_{\mu\nu} = 8\pi G \left[T_{\mu\nu}^{(\text{fl})} + T_{\mu\nu}^{(\phi)} \right]$$

- Non-minimally coupled models

$$S = \int d^4x \sqrt{-g} \left(\frac{1}{2} f(\phi, R) + \mathcal{L}_\phi + \mathcal{L}_\text{fl} \right)$$

$$G_{\mu\nu} = 8\pi G T_{\mu\nu} = \frac{1}{f'} \left[T_{\mu\nu}^{(\text{fl})} + T_{\mu\nu}^{(\phi)} + \frac{1}{2} g_{\mu\nu} (f - f' R) + A_{\mu\nu}(f') \right]$$

$$f(\phi, R) = \frac{F(\phi)}{8\pi G_*} R, \quad F(\phi) = 1 + 8\pi G_* \xi (\phi^2 - \phi_0^2)$$



Basic equations

Continuity equation

$$\dot{\delta} + (1 + \delta)\vec{\nabla}_{\vec{x}} \cdot \vec{u} = 0$$

Euler equation

$$\dot{\vec{u}} + 2H\vec{u} + (\vec{u} \cdot \nabla_{\vec{x}})\vec{u} + \frac{1}{a^2}\nabla_{\vec{x}}\psi = 0$$

Poisson equation

$$\vec{\nabla}_{\vec{x}}^2 \phi_E = \frac{4\pi G}{F} a^2 \bar{\rho}_m \delta$$

$$\phi_E = \left(1 + \frac{1}{2} \frac{F_{,\phi}^2}{F + F_{,\phi}^2} \right) \phi$$

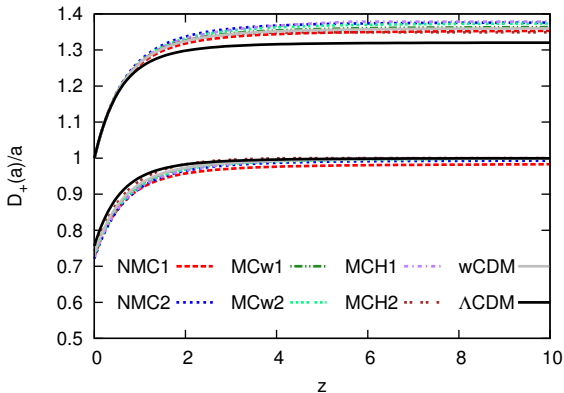
Equations of motion (DM only)

$$\ddot{\delta} + 2H\dot{\delta} - \frac{4}{3} \frac{\dot{\delta}^2}{1 + \delta} - 4\pi G_{\text{eff}} \bar{\rho}_m \delta (1 + \delta) = 0$$

$$\ddot{\delta} + 2H\dot{\delta} - 4\pi G_{\text{eff}} \bar{\rho}_m \delta = 0$$

$$G_{\text{eff}} = \frac{G}{F} \frac{2(F + F_{,\phi}^2)}{2F + 3F_{,\phi}^2} \approx \frac{G}{F} \text{ for } \xi \ll 1$$

Growth Factor

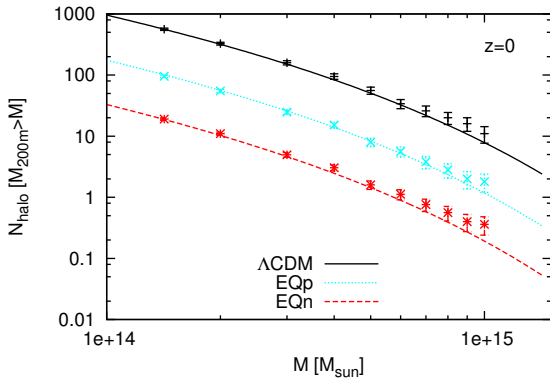


FP et al., 2013

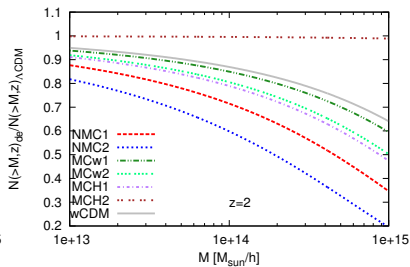
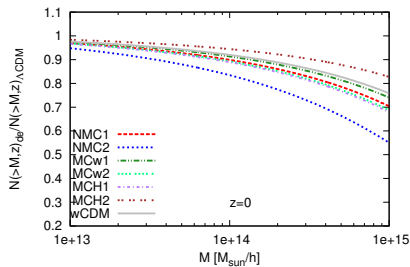
Mass function

- Depends on δ_c , $D_+(a)$ and σ_8
- Sensitive to cosmology in the high-mass tail
- Different models
- Here used the Sheth & Tormen formulation

Comparison with simulations



FP et al., 2013



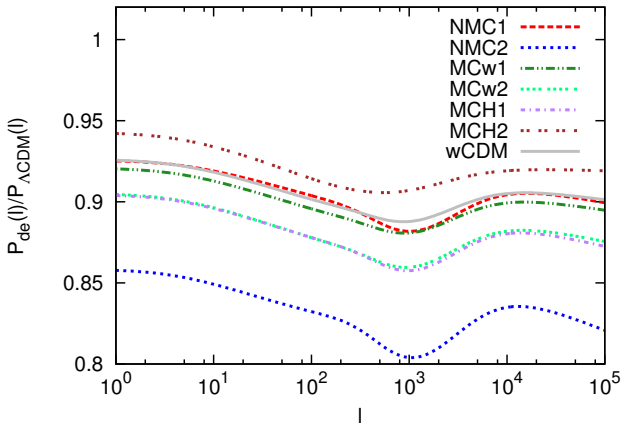
FP et al., 2013

Weak cosmological lensing

- Images of background objects distorted by gravity
- Effect of few percent
- Effective convergence and shear related to matter distribution

$$P_{\kappa}(\ell) = \frac{9H_0^4\Omega_{m,0}^2}{4c^4} \int_0^{\chi_H} \frac{W^2(\chi)}{a^2(\chi)F^2(a)} P_{\delta_m} \left[\frac{\ell}{f_{\kappa}(\chi)}, \chi \right] d\chi$$

Convergence power spectrum



Pace et al., 2013

Conclusions

- Differences in G of order 2% due to solar system constraints
- Small coupling constant
- Appreciable differences in all observables
- Influence of coupling on cluster concentration
- Possibility to probe these models with future surveys