## Cold planar horizons are floppy

## Jorge E. Santos

New frontiers in dynamical gravity


In collaboration with
Sean A. Hartnoll - arXiv:1402.0872 and arXiv:1403.4612

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- What I am going to describe doesn't happen in such setups.

1 The Einstein-Maxwell system

2 Breakdown of Perturbation theory

3 Zero Temperature Numerics

4 Results

5 What about $\mathrm{AdS}_{4}$ ?

6 Conclusion \& Outlook

The bulk theory we study is governed by the Lagrangian

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S=\frac{1}{16 \pi G_{d}} \int \mathrm{~d}^{d} x \sqrt{-g}\left[R+\frac{(d-1)(d-2)}{L^{2}}-\frac{1}{2} F^{a b} F_{a b}\right]
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where $F=\mathrm{d} A$ and $L$ is the $\mathrm{AdS}_{d}$ length scale.

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- Focus on $d=4$, with $\mu(x)=\bar{\mu}\left[1+A_{0} \cos \left(k_{L} x\right)\right]$.
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with $G(y)=1+2 y+3 y^{2}$ and $\delta A_{t}(0, x)=L \sqrt{6} A_{0} \cos \left(k_{L} x\right)$.

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- Brings line element (1) to

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- Solve for the Kodama-Ishibashi variable:

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\begin{aligned}
& \Phi_{-}^{(1)}(\rho, x)=\tilde{\gamma} \cos \left(k_{L} x\right) \rho^{\nu_{-}\left(k_{L}\right)} \text { where } \\
& \qquad \nu_{-}\left(k_{L}\right)=\sqrt{\left(\frac{1}{2}-\sqrt{\frac{k_{L}^{2}}{3}+1}\right)^{2}-\frac{k_{L}^{2}}{6}}-\frac{1}{2}>0
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Breakdown of perturbation theory - resumm perturbation theory.

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- Close to $x=0$, perturbation theory is saved, however away from $x=0$ perturbation theory breaks down!

How to decide which is which?

## Proceed without any approximation - Numerics.

## Ansetzen \& Numerics

Most general line element, without any gauge choice, and compatible with our symmetries takes the following form

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- Use De-Turck method - thank you Toby!


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A=L \sqrt{6}(1-y) P \mathrm{~d} t
\end{gathered}
$$

where $G(y)=1+2 y+3 y^{2}$. For $A=B=S_{1}=S_{2}=P=1$ and $F=0$ it reduces to extreme RN black hole.

## Comments:

- Small irrational powers - $(1-y)^{\nu_{-}\left(k_{L}\right)}-\nu_{-}(1) \approx 0.012$.
- Use finite difference patch near $\mathcal{H}$ and spectral collocation.
- Very steep gradients - need to use adaptive mesh refinement in finite difference patch.
- Use De-Turck method - thank you Toby!
- Alternatively, use very, very small $T / \bar{\mu}$.

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\varpi \equiv \frac{\mathcal{W}_{\max }}{\mathcal{W}_{\min }}-1,
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Einstein's equations chose a resummation that renders the IR floppy - broken translational invariance.

## Emergent picture:



## Periodic potentials in $\mathrm{AdS}_{4}$

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in which case:

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\langle\Phi\rangle_{R}=0, \quad \text { and } \quad\left\langle\Phi_{s}(x, w, 0) \Phi_{s}(s, h, 0)\right\rangle_{R}=\bar{V}^{2} \delta(x-s) \delta(w-h) .
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- What about more general deformations?
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## Outlook:

- Can these new IR geometries affect time dependence?
- Can we make a connection with glassy physics?
- . .

