Cold planar horizons are floppy

Jorge E. Santos New frontiers in dynamical gravity



In collaboration with Sean A. Hartnoll - arXiv:1402.0872 and arXiv:1403.4612

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- What I am going to describe doesn't happen in such setups.



- 2 Breakdown of Perturbation theory
- 3 Zero Temperature Numerics
- 4 Results
- **5** What about AdS₄?
- 6 Conclusion & Outlook

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- Focus on d = 4, with $\mu(x) = \overline{\mu} \left[1 + A_0 \cos(k_L x) \right]$.
- Moduli space space of solutions is 2D: A_0 and $k_0 \equiv k_L/\bar{\mu}$.

• Study time independent perturbations of extremal RN:

$$ds^{2} = \frac{L^{2}}{y^{2}} \left[-G(y)(1-y)^{2}dt^{2} + \frac{dy^{2}}{G(y)(1-y)^{2}} + dx^{2} + dw^{2} \right],$$
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$$\begin{split} \mathrm{d}s^2 &= L^2 \left[\frac{1}{6} \left(-\rho^2 \mathrm{d}\tau^2 + \frac{\mathrm{d}\rho^2}{\rho^2} \right) + \mathrm{d}x^2 + \mathrm{d}w^2 \right] \,, \\ A &= \frac{L\,\rho}{\sqrt{6}} \mathrm{d}\tau \,. \end{split}$$

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• Solve for the Kodama-Ishibashi variable:

$$\begin{split} \Phi^{(1)}_{-}(\rho,x) &= \tilde{\gamma}\cos(k_L x)\rho^{\nu_{-}(k_L)} \quad \text{where} \\ \nu_{-}(k_L) &= \sqrt{\left(\frac{1}{2} - \sqrt{\frac{k_L^2}{3} + 1}\right)^2 - \frac{k_L^2}{6}} - \frac{1}{2} > 0 \end{split}$$

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• Third order Kodama-Ishibashi variable grows faster than first order:

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A tale of two resummations:

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• Close to x = 0, perturbation theory is saved, however away from x = 0 perturbation theory breaks down!

Breakdown of Perturbation theory

How to decide which is which?

Proceed without any approximation - Numerics.

Most general line element, without any gauge choice, and compatible with our symmetries takes the following form

$$ds^{2} = \frac{L^{2}}{y^{2}} \left[-(1-y)^{2} G(y) A dt^{2} + \frac{B}{(1-y)^{2} G(y)} (dy + F dx)^{2} + S_{1} dx^{2} + S_{2} dw^{2} \right]$$
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- Alternatively, use very, very small $T/\bar{\mu}$.



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- For small A_0 , $\varpi \propto A_0$ broken translational invariance $\forall_{A_0 \neq 0} !!!$
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Einstein's equations chose a resummation that renders the IR floppy - broken translational invariance.



Periodic potentials in AdS_4

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in which case:

$$\langle \Phi \rangle_R = 0 \,, \quad \text{and} \quad \langle \Phi_s(x,w,0) \Phi_s(s,h,0) \rangle_R = \bar{V}^2 \delta(x-s) \delta(w-h) \,.$$

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 $\langle g_{ab} \rangle_R$ is accurately described by a Lifshitz geometry:

$$\langle \mathrm{d}s^2 \rangle_R = \frac{L^2}{y^2} \left[-\frac{\mathrm{d}t^2}{y^{2(\bar{z}-1)}} + \mathrm{d}x^2 + \mathrm{d}w^2 + \mathrm{d}y^2 \right] \,.$$

Conclusions:

- Numerical evidence that $\mathsf{AdS}_2\times\mathbb{R}^2$ is RG unstable.
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Outlook:

- Can these new IR geometries affect time dependence?
- Can we make a connection with glassy physics?
- . . .